A numerical typhoon track prediction model and its application on the real-time operation

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ABSTRACT. Recently a nested numerical model for typhoon track prediction (MTTP) with a higher horizontal resolution and a complex physical package has been developed by the National Meteorological Center (NMC, Beijing). For its improvement of initialization, a modified typhoon bogus scheme (Iwasaki 1987) also has been applied.

The design of the MTTP had been determined by the end of 1992 and the forecast capability of this model was tested firstly with a series of experiments for the selected typhoon cases in 1993 (Wang and Li 1994). After it was examined in real-time forecast in the next year, further improvements were made in 1995 including a higher horizontal resolution increased from about 100 to 50 km, a package of complex physics instead of the simple one and a scheme for removal of analyzed vortex. With the improved forecast capability, the MTTP run in quasi-operation from the date of 1 June 1995. Its products were also used by the forecasters during the past two years and the results were very encouraging.

Key words — Nested numerical model, Physical processes, Typhoon bogus, Track prediction, Operational application.

1. Introduction

In history, the NWP of the tropical cyclone (TC) has been such a difficult work, because it is sensitive to the meteorological observations over the tropical ocean as well as the model resolution and physical process parameterization. During the past thirty years, because of the significant advances on the computer processing speed and available memory and the continuing research progress in tropical meteorology, especially in the TC knowledge, it is now possible for better TC forecasts with the rapid development of high resolution models with the package of more detailed and advanced physics.

Recent theoretical research (Holland 1984, DeMaria 1985) demonstrated that TC motions were strongly influenced by the two main factors: one was the
asymmetric basic current and the other was the TC horizontal scale and its outer structure. However, due to the scarcity of the conventional observation on the large ocean, the TC circulation was poorly produced by the objective analysis, with its inaccurate location. So the initial condition formed by the regular observational data failed to resolve or simulate the above important factors. This real requirement prompted the development of TC bogus technique.

Based on the above advantage and funds by the project 85-906 of the State Science Committee of China, the (MTTP) has been implemented at NMC from the development of the limited area analysis and forecast system (LAFFS) (Guo et al. 1992). Especially focusing on the typhoon track prediction, the model is a nested mesh-grid one with a higher horizontal resolution (about 50 km) and a package of complex physics. For the improvement of the model initial condition, a typhoon initialization procedure was made with the recent advanced knowledge on (i) removal of the analyzed typhoon vortex and (ii) insertion of the bogus TC into original analysis field. The MTTP performed considerably forecasts skillful for typhoon track by using the above schemes.

2. General specification of the MTTP

2.1. Governing equations

The governing equations are composed of the momentum, mass and moist continuity, thermodynamic and hydrostatic equations, written in the flux form as follows:

\[
\frac{\partial u}{\partial t} \cos \theta \left[ \frac{1}{P_s} \frac{\partial T}{\partial \theta} + P_s \nu \cos \theta \frac{\partial T}{\partial \theta} \right] + \frac{1}{a \cos \theta} \left[ \frac{\partial}{\partial \theta} (P_s u) + P_s \nu \cos \theta \frac{\partial u}{\partial \theta} \right] + \frac{\partial}{\partial \sigma} (P_s \sigma) = 0
\]

\[\frac{\partial v}{\partial t} + ZP_s \nu \frac{\partial}{\partial \theta} (P_s \nu) + RT \frac{\partial}{\partial \theta} \ln P_s + \hat{\sigma} \frac{\partial u}{\partial \sigma} = F_u + D_u\]

\[\frac{\partial P_s}{\partial t} + \frac{1}{P_s} \frac{\partial}{\partial \theta} \left[ \frac{\partial}{\partial \theta} (P_s \mu) + \frac{\partial}{\partial \theta} (P_s \nu \cos \theta) \right] + \frac{\partial}{\partial \sigma} (P_s \sigma) = 0\]

\[\frac{\partial T}{\partial t} + \frac{1}{P_s} \frac{\partial}{\partial \theta} \left[ \frac{\partial}{\partial \theta} (P_s \mu) + P_s \nu \cos \theta \frac{\partial T}{\partial \theta} \right] + \frac{\partial}{\partial \sigma} (P_s \sigma) = Q + D_T\]

\[\frac{\partial \hat{\sigma}}{\partial \ln \sigma} = -RT\]

where, \(Z\) is the absolute potential vorticity and \(E\) is the unit-mass kinetic energy given in detail below:

\[Z = \frac{1}{P_s} \left\{ f + \frac{1}{a \cos \theta} \left[ \frac{\partial v}{\partial u} - \frac{\partial}{\partial \theta} (u \cos \theta) \right] \right\}\]

\[E = \frac{1}{2} \left( u^2 + \frac{1}{\cos \theta} v^2 \cos \theta \right)\]

where, \(\alpha = R/C_p\), geostrophic parameter, \(f = 2\Omega \cos \theta\) and \(\omega\) is the vertical velocity in the pressure coordinate defined as,

\[\omega = P_s \frac{\partial}{\partial t} + P_s \nu \cos \theta \frac{\partial}{\partial \theta} \ln P_s\]

In the above equations, \(Q\) includes the diabatic heating arising from radiation and latent heat release, \(S\) includes the moisture source/sink arising from cumulus convection and large-scale condensation, \(F_u\) and \(F_v\) are surface frictional force terms and \(D_u\), \(D_v\), \(D_q\), \(D_t\) denote the horizontal diffusion terms. The Eqsns. (1)-(9) use the boundary condition as the following:

\[(P_s \hat{\sigma})_{\sigma=0} = 0, (P_s \hat{\sigma})_{\sigma=1} = 0\]

2.2. Grid point generations

The sigma coordinate system is utilized in the vertical with 15 irregular levels, defined as

\[\sigma_k = 0.75 \left( \frac{k}{15} \right) + 1.75 \left( \frac{k}{15} \right)^3 - 1.5 \left( \frac{k}{15} \right)^4\]

\[(k = 1, 2, 3, ..., 15)\]

where, \(\sigma = P/P_s\), \(P\) is the pressure and the subscript \(s\) denotes the surface value. The vertical distribution of variables is that the horizontal wind \((U, V)\), temperature \((T)\) and specific humidity \((q)\) are on the integer levels while the vertical velocity \((\hat{\sigma})\) and geopotential height \((\hat{\phi})\) are on the semi-levels.
The model integrations are performed on the staggered Arakawa C grid using the reaped difference for the dynamic and forward difference for the diabatic heating.

2.3. Lateral boundary

A one-way boundary scheme is used between the fine mesh and the coarse as well as between the coarse mesh and the global model (T63L16). At the four outermost rows of fine-mesh gridpoints, Davies relaxation method is applied to specify the boundary value as

\[ S_{\tau+1}^* = (1-\alpha) S_{\tau+1}^* + \alpha S_{\tau+1} \]  

(12)

where \( S_{\tau+1}^* \) is any fine-mesh model available while \( S_{\tau+1}^* \) is the original one before this process. \( S_{\tau+1} \) is the corresponding value of the coarse-mesh,

\[ S_{\tau+1} = S_{\tau} + \Delta S \]  

(13)

where \( S_{\tau} \) is any fine-mesh value at \( \tau \) time and \( \Delta S \) is the boundary tendency from the coarse-mesh model. Here \( \alpha \) represents the weighting function as follows,

\[ \alpha = 1 - \tanh (kk) \]  

(14)

\( kk \) is the specific number of fine-mesh point from the boundary, \( kk = 1, 2, 3, 4 \).

2.4. Forecast domain

The MTTCP is a regional multiple nested model. Its first guess fields are supplied by the operational global spectral model system (T63L16) of NMC. The horizontal resolution of the coarse mesh is 0.9375° over the domain 0°N-50.625°N, 84.375°E-161.25°E, while that of the fine mesh is 0.46875° (approximately 50 km) over the domain 10.3125°N-41.25°N, 105°E-150.9375°E. The grid structure, controlling equations, lateral boundary, time integration and physical processes schemes are all the same for both the coarse and the fine mesh. The forecast domain could be extended if the condition of computer sources is allowed.

3. Physical processes

In the last three decades, many studies have been done and the significant progress had been attained in the field of physical processes, such as, the work by Kuo (1974), Anthes (1977), Louis (1979), Tibaldi (1982), Tiedtke (1984), Hart et al. (1988), Puri et al. (1992) and so on. Following their advances, a package of complex physics is implemented in both the coarse-mesh and fine-mesh grid of the model. Here only a brief description will be given.

3.1. Deep cumulus convection parameterization

A new version of Kuo (1974) scheme is used for the vertical integration treatment of the local moisture accumulation: (i) the moisture for the occurrence and development of cumulus convection is supplied by the large-scale environment but not the large-scale moisture convergence; (ii) the multiple cumulus layers are allowed to exist simultaneously.

3.2. Shallow cumulus convection parameterization

Using the scheme by Tiedtke (1987), the treatment of shallow convection is similar to that of the turbulent diffusion. The effect produced by the shallow convection does not enhance the heating by liquid precipitation, but only adjusts the profile of atmospheric heat and water vapour. The occurrence of shallow convection is specified according to the atmospheric conditional instability. When the layer for shallow convection is determined, it will be used to solve the simultaneous equations with the turbulent flux forms used in the vertical diffusion equation.

3.3. Vertical turbulent diffusion parameterization

The treatment of vertical diffusion is similar to that of turbulent transport in the atmospheric boundary layer (K theory). The subgrid-scale processes of vertical diffusion of heat, momentum and moisture are expressed as the forms of diffusion flux so that their parameterization can be specified by determination of the coefficients (K) of diffusion and exchange. The coefficients are calculated with the formulas by Louis (1979), which is relative to the stability. For any variable \( \phi \) vertical diffusion equation using implicit difference scheme can be written as,

\[ \left( \frac{\partial \phi}{\partial t} \right) = \frac{1}{\rho} \frac{\partial}{\partial z} \left( \bar{\rho} K_\phi \frac{\partial \phi}{\partial z} \right) \]  

(15)

giving the upper boundary condition for the above formula: when \( P = Pt \), the right hand term is set to zero; and the lower boundary condition: when \( P = Ps \), the right hand term is set to \( C_\phi \left| \nabla \right| (\phi_1 - \phi_s) \), where the subscripts 1 and s indicate the lowest model level and the surface
respectively. $C_d$ is the drag coefficient of surface turbulent flux that varies depending on the different value of the Richardson parameter (RI = 0, neutral; RI > 0, stability; RI < 0, instability). $K_g$ is defined as the turbulent transfer coefficient that is the similarity function of RI and the model layer.

3.4. Radiation processes

The schemes of Lacis and Hansen (1974) and Fels and Schwarzkopf (1975) are introduced here for the radiative transfer calculations of solar wavelengths and terrestrial wavelengths. In order to reduce the CPU time, the tables are used for the calculations of radiation radio. Cloud amounts and heights are fixed zonal averages for each season using climatological data.

3.5. Surface thermal budget

The total surface budget can be determined by the terms defined as the following: the surface absorption of solar short-wave and atmospheric longwave radiation ($S_g$ and $A_g$ respectively), the terrestrial surface radiation ($\sigma_g T_s^4$), the conduction of surface latent heat [$-H_g(T_s)$] and sensible heat [$-H_g(T_s)$] to the atmosphere, the conduction of surface heat to the subsurface soil or ocean [$-H_g(T_s)$] and the heat transfer specified by the evaporation at the ice surface [$H_{ice}(T_s)$]. So the heat balance equation at the surface is given by:

$$S_g + A_g - \sigma_g T_s^4 - H_g(T_s) - H_g(T_s) - H_g(T_s) - H_{ice}(T_s) = 0$$  \hspace{1cm} (16)

where $\sigma_g$ is the Stefan-Boltzmann constant. An iteration method may be used to solve the equation so as to obtain the surface temperature ($T_s$) while the influence by the different condition of underlying surface is also considered in the scheme.

3.6. Soil thermal conductivity

Three subsurface layers at depths of 0.05, 0.5 and 5 m is used to calculate the subsurface temperature. The conduction equation of soil heat has the form

$$\frac{\partial T_s}{\partial t} = \frac{\partial}{\partial z} \left( K_s \frac{\partial T_s}{\partial z} \right)$$  \hspace{1cm} (17)

where giving the assumption: $K_s(0) = K_s(R_z) T_s = 280^\circ$. $K_s$ denotes the thermal diffusivity and $T_s$ is the ground temperature. All of the surface coefficients used in the calculation, such as reflectivity, roughness, surface temperature, soil moisture, snow depth and SST etc. are fixed to the monthly climatological data from the global model (T63L16).

3.7. Large-scale condensation

The saturated condensation scheme (Guo et al. 1992) is applied as the following: when the over-saturation occurs in the air of model, the rainout are made to return the relative humidity to 100%. So the existing of condensed cloud is not permitted in this scheme. Besides that, the process that the liquid water evaporates again during it descends is also considered.

3.8. Horizontal diffusion parameterization

A non-linear second order diffusion scheme is used as follows.

$$D_k = \kappa \left| \nabla^2 \chi \right| \nabla^2 \chi$$  \hspace{1cm} (18)

where, $\chi$ is any variable, $\kappa$ is the diffusion coefficient and $\nabla^2$ represents Laplace operator as follows:

$$\nabla^2 \chi = \frac{\Delta^2 \chi}{a^2 \cos^2 \theta \Delta \lambda^2} + \frac{\Delta (\cos \theta \Delta \chi)}{a^2 \cos \theta \Delta \theta^2}$$

3.9. Moisture modification

Because the advection term at the moisture equation is calculated using the non-regular difference method, the incorrectness that the negative moisture occurs can not be avoided. So the modification must be made for those points where moisture are negative. Such a negative value of the central point is compensated to zero by those of the four surrounding points on the vertical meridional cross-section. The contribution from a surrounding point is specified by its value and the average weight determined from its distance to the compensated point. A shorter distance and a larger moisture value may lead to a larger weight coefficient while the opposite occurs, if the contrary situation prevails.

4. Removal of the analyzed typhoon vortex from the objective analysis field

Due to the lack of observation in the ocean, the current analysis has a difficulty in describing the detailed structure of typhoon vortex. The analyzed vortices were evidently too large and weak with their inaccurate central positions and intensities that were different from their corressonspible realistic ones. Such
a initial condition for model would result in the original forecast errors that significantly affect the typhoon motion. So the poorly resolved analyzed vortex should be removed from the large-scale analysis before the bogus typhoon is blended into the analysis field.

The scheme of Kurihara et al. (1990, 1993) is introduced for solving the problem. Here, the original analysis field \( h \) can be split into the large-scale basic field \( (h_B) \) and the disturbance field \( (h_D) \) as follows,

\[
h = h_B + h_D
\]

(19)

where, \( h \) represents any variable. Let the analyzed typhoon vortex \( (h_{av}) \) be included in the disturbance \( (h_D) \), then the environmental field \( (h_E) \) can be obtained by recombining the remaining non-typhoon disturbance and the basic field as :

\[
h_E = h_B + (h_D - h_{av})
\]

(20)

or \( h_E = h - h_{av} \)

(21)

The filtering is carried out in both the zonal and meridional directions respectively to divide \( h \) into \( h_D \) and \( h_B \), then \( h_{av} \) can be separated from \( h_D \) using the cylindrical filtering. In detail, filtering in the zonal direction is firstly applied using a local three-point smoothing operator as follows,

\[
\tilde{h}_{i,j} = h_{i,j} + k(h_{i-1,j} + h_{i+1,j} - 2h_{i,j})
\]

(22)

where, \( \tilde{h} \) is the zonal-smoothed value and the subscript \( i \) and \( j \) denote longitude and latitude respectively, \( k \) is the filtering coefficient defined by,

\[
k = \frac{1}{2} \left( 1 - \cos \frac{2\pi}{m} \right)^{-1}
\]

(23)

where, \( m \) sequentially varies as 2, 3, 4, 2, 5, 6, 7, 2, 8, 9, 2 so the components with less than 9 degrees wavelength will be filtered. The above procedure is applied for the filtering in the meridional direction to obtain the basic field as the following:

\[
\tilde{h}_{i,j} = \tilde{h}_{i,j} + k(\tilde{h}_{i-1,j} + \tilde{h}_{i+1,j} - 2\tilde{h}_{i,j})
\]

(24)

Then the \( h_D \) can be easily obtained from the Eqn. (19).

The isolation of the analyzed typhoon vortex \( h_{av} \) from \( h_0 \) is accomplished through the cylindrical filtering in the polar coordinate system \( (r, \theta) \) written by,

\[
h_{av}(r, \theta) = [1 - E(r)] [h_D(r, \theta) - \tilde{h}_{av}(r_0)]
\]

(25)

where, \( E(r) \) is the empirical function to determine the filtering characteristic that \( h_{av} \) gradually diminishes to zero at the filter edge as the following:

\[
E(r) = \begin{cases} 
\exp \left\{ -\frac{(r_0 - r)^2}{l^2} \right\} - \exp \left\{ -\frac{(r_0)^2}{l^2} \right\} & r \leq r_0 \\
1 & r > r_0 
\end{cases}
\]

(26)

In this formula, \( l \) controls the filtering shapes and is empirically set to \( r_0/5 \) here. \( r_0 \) is the radius from the center of analyzed typhoon vortex and defines the filtering domain.

The \( r_0 \) determination and the center of analyzed typhoon vortex are described below.

Let \( V_D \) be the disturbance wind velocity at the 850 hPa level given by,

\[
V_D = \left( \frac{\nabla D + V_D}{2} \right)^{1/2}
\]

(27)

Then the central position \( (\lambda_0, \varphi_0) \) of the analyzed typhoon vortex can be specified by the distribution of \( V_D \) as follows,

\[
\lambda_0 = \frac{\sum V_{Di,j} \lambda_{i,j} \Delta S_{i,j}}{\sum V_{Di,j} \Delta S_{i,j}}
\]

(28)

\[
\varphi_0 = \frac{\sum V_{Di,j} \varphi_{i,j} \Delta S_{i,j}}{\sum V_{Di,j} \Delta S_{i,j}}
\]

(29)

where, \( i \) and \( j \) indicate the longitude and latitude respectively and \( \Delta S_{i,j} \) is the area of the assigned grid point. In the current analysis field, the disturbance wind is strong enough in the analyzed typhoon vortex region so the disturbance radius \( r_0 \) can be determined from the profile of \( V_D(r) \) that defined the circular mean of \( V_D \) at a radius \( r \) from the center \( (\lambda_0, \varphi_0) \). Empirically the \( V_D(r) \) decreases with radius after it attains a maximum and reaches a relatively small value level at which the radius \( r_0 \) can be specified. Alternatively, the \( r_0 \) would be set to a limit if above profile could not be met. Then the average along the periphery of \( h_D \) at the radius
$r_0$ can be calculated as follows:

$$\bar{h}_D (r_0) = \frac{1}{2\pi} \int_0^\Phi h_D (r_0, \theta) d\theta$$  \hspace{1cm} (30)

By Eqn. (25), an axisymmetric analyzed typhoon vortex will be finally isolated from the analysis field. Its intensity decreases smoothly from the center to zero at the filtering edge.

Then the $h_E$ can be calculated from the Eqn. (21) and will be used for the environmental field where the bogus typhoon be set up.

5. Typhoon bogus

The empirical method is applied to use the available surface parameters, such as the observed typhoon positions at current time and previous 12 hours, central pressure, radius of 15 m/s winds etc., to construct a bogus typhoon and blend it with the large-scale environmental field as the composite initial condition for the model forecast.

Following the work by Iwasaki (1987), the revised typhoon bogus scheme is described briefly as below.

5.1. Surface pressure field

The Fujita formula (1952) is used to define a symmetric surface pressure distribution as follows:

$$p(r) = p_E - \Delta p \left[ 1 + \left( r/R_0 \right)^2 \right]^{1/2}$$  \hspace{1cm} (31)

where, $p_E$ is the mean environmental pressure, $\Delta p$ is the difference between $p_E$ and the observed central pressure. $R_0$ is the factor to determine the profile of the radial pressure gradients around the typhoon vortex center. It is decided by the radius of 15 m/s winds ($R_{15}$) and the Eqn. (31). Here $R_0$ is limited to no less than 30 km.

5.2. Upper anticyclone

Above the cloud top level ($p_t$), it is assumed that an anticyclone exists and the $D$-value is used to describe the deviation between typhoon and environmental field. The $D$-value is defined as,

$$D(r, p_t) = \begin{cases} 
  ar^2 + b & r > R_0 \\
  cr + d & R_0 < r < R_1 \\
  e(r - R_E)^2 & R_1 < r < R_E 
\end{cases}$$  \hspace{1cm} (32)

where, $R_1$ and $R_E$ are proportional to $R_{15}$. The parameters $b$ and $\Delta p$ represent the strength of the upper anticyclone at $p_t$ level at which the environmental temperature sounding is intersected with the wet adiabatic profile rising from the sea surface. The parameters $a$, $c$, $d$ and $e$ can be calculated by the conditions that keep $D$-value profile changing smoothly at the points $R_1$ and $R_E$. Let $D$-value and its first-order difference be equal at the two points, the following formula can be derived as follows:

$$\begin{align*}
  aR_0^2 + b &= cR_0 + d \\
  cR_1 + d &= e(R_1 - R_E)^2 \\
  2aR_0 &= c \\
  c &= 2e(R_1 - R_E)
\end{align*}$$  \hspace{1cm} (33)

Using the given $R_0, R_1, R_E$ and $b$ values, the parameters $a$, $c$, $d$ and $e$ are computed using the above formula and the anticyclone will be determined with assumption that the $D$-value vanish at 20 hPa ($p_{mid}$), i.e.,

$$D(r, p_{mid}) = 0$$  \hspace{1cm} (34)

5.3. Typhoon warm core

The vertical profile of the temperature at bogus typhoon center below the cloud top level ($p_t$) is prescribed by,

$$T(0, p) = C_1 [T_c(p) - T_E(p)] + T_E(p)$$  \hspace{1cm} (35)

where, $T_E$ is the environmental temperature. $T_c$ is the moist adiabatic temperature rising from sea surface temperature. The constant $C_1$ is determined by the $D$-value from the deviation between $T_c$ and $T_E$ at $p_t$ level according to the static equilibrium and the Eqn. (32).

Above the cloud top level ($p_t$), assuming the temperature at typhoon center, vary with logarithmic pressure along parabola, it is given by

$$T(0, p) = C_2 (\ln p - \ln p_{mid}) (\ln p - \ln p_t) + T_E(p)$$  \hspace{1cm} (36)

where, the constant $C_2$ is defined from Eqns. (31) and (33) and the static relationship. Then the $D$-value at typhoon center at the pressure layers can be determined from the temperature deviation between
\( T_c \) and \( T_E \).

At other radii, the distributions of \( D \)-value are specified from the formula:

\[
D(r, p) = \alpha(r) D(0, p) + \beta(r) \tag{37}
\]

The coefficients \( \alpha \) and \( \beta \) are determined by Eqns. (31) and (32) below the cloud top level \( (p_t) \) and by Eqns. (32) and (34) above \( p_t \) then written as,

\[
\alpha(r) = \frac{[D(r, p_t) - D(r, p)]}{[D(0, p_t) - D(0, p_t)]} \quad (p \leq p_t)
\]

\[
\beta(r) = \frac{[D(r, p_t) D(0, p_t) - D(0, p_t) D(r, p_t)]}{[D(0, p_t) - D(0, p_t)]}
\]

and

\[
\alpha(r) = \frac{D(r, p_t)}{D(0, p_t)} \quad (p \geq p_t)
\]

\[
\beta(r) = 0
\]

where, \( p_t \) represents the sea surface level where the \( D \)-value is calculated from Eqn. (31) and the static relationship. Moreover, the moisture distribution is given by 90% relative humidity around the typhoon below the cloud top level \( (p_t) \).

5.4. Gradient wind field

The maintenance of typhoon circulation depends on the convergence at the lower layer. So the steady momentum equations neglecting the vertical advection are used to calculate the gradient wind. Including surface friction this equation in the cylindrical coordinate can be reduced as,

\[
v_r \frac{\partial v_r}{\partial r} - \frac{v_\theta^2}{r} - f v_\theta + \frac{\partial \phi}{\partial r} + C_d l v \nu_r = 0
\]

\[
v_r \frac{\partial v_\theta}{\partial r} + v_r v_\theta + f v_r + C_d l v \nu_\theta = 0
\] \tag{38}

where, \( v_r \) and \( v_\theta \) are radial wind and tangential wind component respectively, \( f \) is the Coriolis parameter, \( C_d \) represents the drag coefficient and \( \phi \) is the geopotential height. In the upper troposphere, assumption is made that the divergence be set to compensate the convergence at the lower layers.

5.5. Typhoon Initial motion

The initial movement vector is determined from the positions with a 12 hours interval. This vector is then uniformly added to the wind field of the bogus typhoon. The compensation for the mass field associated with the initial vector is neglected here.

5.6. Implantation of bogus typhoon in the environment

The bogus typhoon vortex is superposed on to the original analysis field using the weighted mean technique as follows,

\[
f(r, \theta) = f_m (r, \theta) w(r) + f_E (r, \theta) [1 - w(r)] \tag{39}
\]

where, \( f_m \) and \( f_E \) are bogus typhoon and analysis field respectively, the radius and azimuth angle \( (r, \theta) \) denote the grid points in polar coordinate. The weight coefficient \( w(r) \) is defined as,

\[
w(r) = \begin{cases} 
1 & r \leq R_{inner} \\
\cos \left( \frac{\pi}{2} \cdot \frac{r - R_{inner}}{R_{outer} - R_{inner}} \right) & R_{inner} \leq r \leq R_{outer} \\
0 & R_{outer} < r
\end{cases}
\] \tag{40}

where \( R_{inner} \) and \( R_{outer} \) are proportional to \( R_{15} \).

6. Operational results

The first version of MTTP had been completed with the simple physical processes and a lower horizontal resolution (approximately 100 km) by the end of 1992. In the next year, 26 forecasts were made for 10 typhoons or TCs locating in the model domain, which the prediction capabilities of MTTP were initially tested and the results showed feasible. In 1994, total 89 real-time forecasts for 23 typhoons or TCs were produced by MTTP for its further experiments and some of them were used as reference in operation. The results showed that the model was encouragingly skillful in forecasting typhoon tracks but poor in maintaining their central pressures and had somewhat trouble in the false spin-up in the initial time (not presented here). The problem, most possibly, was mainly due to the lower horizontal resolution, the simpler physics package and the incomplete initialization scheme.

Based on the above experiments, the main improvements were made for the MTTP in 1995 including the introduction of complex physics package instead of the simple one, a higher horizontal resolution
Fig. 1. The observed and predicted tracks of typhoon Faye (T9503)

Fig. 2. The observed and predicted tracks of severe tropical storm Gray (T9504)

<table>
<thead>
<tr>
<th>Year</th>
<th>t = 24 hours</th>
<th>t = 48 hours</th>
</tr>
</thead>
<tbody>
<tr>
<td>1993</td>
<td>167 (26)</td>
<td>382 (23)</td>
</tr>
<tr>
<td>1994</td>
<td>196 (89)</td>
<td>361 (61)</td>
</tr>
<tr>
<td>1995</td>
<td>183 (85)</td>
<td>332 (69)</td>
</tr>
</tbody>
</table>

*The numbers in parentheses denote the forecast cases (about 50 km) upgraded from the coarser one (about 100 km) and a new typhoon initialization scheme. The revised version of MTTP was installed and upgraded in quasi-operation from 1 June 1995. Then MTTP run for 85 forecasts of 17 typhoons or TCs and its products were used in the real-time operation for typhoon track prediction by NMC (Beijing). Its forecast track error statistic and those during the prior two years are presented as Table 1.

Some forecast cases of the typical typhoons or TCs that affected the coast of China are presented in detail below and shown in Figs. 1-6, where the character Fs denote the forecast central positions of typhoon or TC while the character Zs are the corresponding observed ones.

6.1. The events of recurvature: typhoon Faye (T9503), tropical severe storm Gary (T9504) and tropical severe storm Helen (T9505)

T9503 was a case of northward recurvature typhoon that made landfall on the South Korean coast on the mid-night of 23 July 1995. Total six day forecasts for the case were made from 18 July and the MTTP simulated the typhoon’s northward tendency realistically. Among them, the forecasts on the two days of 18 and 20 July were better than those on 19 and 21 July. The best one is the forecast on 22 July prior to the date of landfall, which the model succeeded in predicting the landfall location and the landfall time (Fig. 1).

T9504 is also a recurvature case and the TC finally landed on the area between Chenghai and Yaopin in Guangdong Province of China. The TC moved west since it formed on 0600 UTC 29 July 1995. Generally it was estimated that it would keep its motion westward, but in fact it turned north rapidly from the mid-night of 30 July due to the influence by the development of South-Sea anticyclone. This recurvature process was simulated successfully by the model based on the data of 0000 UTC 30 July (Fig. 2). The encouraging forecasts for T9503 and T9504 brought the MTTP to the attention of the forecasters at NMC (Beijing).

The case T9505 is similar to T9504. Formed on 1200 UTC 9 August 1995, the TC shifted westward then rapidly turned north and made landfall on the area between Shenzhen and Huizhou in Guangdong Province of China. A series of forecasts for the TC were made with model running twice a day. The model was successful in forecasting the TC movement tendency, among which the second and the fourth were better than the first and the third. The forecast
TC position errors for landfall compared with the observation were considerably little (Fig. 3).

6.2. The case with re-entry to sea: tropical severe storm Janis (T9507)

The TC moved northwest after it formed and landed on Shiting town of Wenzin in Zhejiang Province of China in the morning of 25 August 1995. Then it turned northeast and entered Yellow-Sea again to shift to the Korean. The forecasts of MTTP were made from 1200 UTC 22 August and the model simulated the TC movement tendency realistically. Compared with the 48-hour forecasts, the 24-hour ones were better for the TC recurvature and its landfall time and location (Fig. 4).

6.3. The event affected by the island orography: tropical severe storm Lois (T9508)

This TC formed on the South-Sea and then made landfall rapidly on Wangning in Hainan Province of China in the morning of 28 August 1995. The MTTP had a little bad performance in landfall of this case (Fig. 5). Initial diagnosis and other prior studies showed that the effect of island orography on typhoons or TCs had not been expressed well in the current model design (not presented here).
6.4. The case of rapid acceleration: typhoon Kent (T9509)

This typhoon did not recurve but moved slowly on the sea surface on the east of Philippine. From 29 August 1995, it accelerated rapidly and then landed on the coast between Haifong and Huidong in Guangdong Province of China in the afternoon of 31 August. The model had large forecast error in simulating the typhoon's slow shift but had excellent performance in its rapid motion from 29 August. Especially the prediction for landfall location and time was very realistic (Fig. 6).

7. Conclusions

From above MTTP forecast cases and the performance of this model in both 1994 and 1995, it could be indicated that the MTTP had the considerable capability in the prediction of typhoons or TCs whose tracks were normal, and somewhat skills for other abnormal tracks such as those turning or accelerating rapidly. These results showed that the MTTP had a position in the real-time operation with its useful forecast production. However, the MTTP still had inability in some aspects and its improvement will be continued in the coming future.

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References


